

NNLL resummation for QCD cross sections

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I present results for the resummation of soft-gluon contributions to QCD hard-scattering cross sections at next-to-next-to-leading logarithm accuracy. A key ingredient is the calculation of two-loop soft anomalous dimensions for the partonic processes. Explicit expressions and applications are provided for processes that involve massless partons and/or massive quarks.

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1. Resummation

QCD corrections can be quite large for many processes in hadron colliders. Certain regions of kinematical phase space can contribute the dominant part to these corrections and can be significant even beyond NLO. One such region is partonic threshold, where soft-gluon corrections are important. These corrections are needed for increased accuracy in theoretical predictions and, using factorization theorems and renormalization group evolution, they can be resummed in terms of soft anomalous dimensions. Two-loop eikonal calculations of these quantities allow the resummation to be performed at next-to-next-to-leading logarithm (NNLL) accuracy.

We resum logarithms of a moment variable N , conjugate to a kinematical variable that measures distance from threshold. We then write the resummed cross section as [1]

$$\sigma(N) = \exp \left[\sum_i E_i(N) \right] H(\alpha_s) \exp \left[\int_{\sqrt{s}}^{\sqrt{s}/\tilde{N}} \frac{d\mu}{\mu} \Gamma_S^\dagger(\mu) \right] S \left(\frac{\sqrt{s}}{\tilde{N}} \right) \exp \left[\int_{\sqrt{s}}^{\sqrt{s}/\tilde{N}} \frac{d\mu}{\mu} \Gamma_S(\mu) \right]$$

where E_i resums universal collinear and soft-gluon emission from the incoming partons, H is the hard-scattering function, S is the soft-gluon function describing noncollinear soft-gluon emission, and Γ_S is the soft anomalous dimension - a matrix in color space and a function of kinematical invariants. We calculate Γ_S in the eikonal approximation, in momentum space and using Feynman gauge, from the UV poles of dimensionally-regularized eikonal diagrams.

2. Two-loop soft anomalous dimensions

Complete two-loop results for Γ_S are now available for the soft (cusp) anomalous dimension for $e^+e^- \rightarrow t\bar{t}$, for s -channel single top production, for $bg \rightarrow tW^-$ and $bg \rightarrow tH^-$, and for $t\bar{t}$ hadroproduction.

2.1 Soft (cusp) anomalous dimension for $e^+e^- \rightarrow t\bar{t}$

The soft (cusp) anomalous dimension for $e^+e^- \rightarrow t\bar{t}$ is $\Gamma_S = (\alpha_s/\pi)\Gamma_S^{(1)} + (\alpha_s/\pi)^2\Gamma_S^{(2)} + \dots$ with $\Gamma_S^{(1)} = C_F(\gamma \coth \gamma - 1)$, in terms of the cusp angle $\gamma = \ln[(1+\beta)/(1-\beta)]$ with $\beta = \sqrt{1-4m^2/s}$ and m the top quark mass, and [2]

$$\Gamma_S^{(2)} = \frac{K}{2}\Gamma_S^{(1)} + C_F C_A \left\{ \frac{1}{2} + \frac{\zeta_2}{2} + \frac{\gamma^2}{2} - \frac{1}{2} \coth^2 \gamma \left[\zeta_3 - \zeta_2 \gamma - \frac{\gamma^3}{3} - \gamma \text{Li}_2(e^{-2\gamma}) - \text{Li}_3(e^{-2\gamma}) \right] \right. \\ \left. - \frac{1}{2} \coth \gamma \left[\zeta_2 + \zeta_2 \gamma + \gamma^2 + \frac{\gamma^3}{3} + 2\gamma \ln(1 - e^{-2\gamma}) - \text{Li}_2(e^{-2\gamma}) \right] \right\}$$

where $K = C_A(67/18 - \zeta_2) - 5n_f/9$.

The cusp anomalous dimension written above is an essential component of calculations for other QCD processes, where the color structure gets more complicated with more than two colored partons involved in the process.

2.2 Soft anomalous dimensions for single-top production processes

Next, we compute two-loop soft anomalous dimensions for single top production in the s -channel and for the associated production of a top quark with a W boson or a charged Higgs.

For s -channel single top production the one-loop [3, 4] and two-loop [4] expressions are

$$\begin{aligned}\Gamma_{S,\text{top } s\text{-ch}}^{(1)} &= C_F \left[\ln \left(\frac{s-m^2}{m\sqrt{s}} \right) - \frac{1}{2} \right] \\ \Gamma_{S,\text{top } s\text{-ch}}^{(2)} &= \frac{K}{2} \Gamma_{S,\text{top } s\text{-ch}}^{(1)} + C_F C_A \frac{(1-\zeta_3)}{4}.\end{aligned}$$

For associated production of a top quark with a W^- , $bg \rightarrow tW^-$, the soft anomalous dimension at one-loop [3, 5] and two-loops [5] is

$$\begin{aligned}\Gamma_{S,tW^-}^{(1)} &= C_F \left[\ln \left(\frac{m^2-t}{m\sqrt{s}} \right) - \frac{1}{2} \right] + \frac{C_A}{2} \ln \left(\frac{m^2-u}{m^2-t} \right) \\ \Gamma_{S,tW^-}^{(2)} &= \frac{K}{2} \Gamma_{S,tW^-}^{(1)} + C_F C_A \frac{(1-\zeta_3)}{4}.\end{aligned}$$

We find the same analytical results for the related process $bg \rightarrow tH^-$ [6, 5].

2.3 Soft anomalous dimension matrices for $t\bar{t}$ production

For top-antitop production in hadron colliders, the soft anomalous dimensions were derived at one-loop in [1] and have been studied in various approaches at two-loops in [2, 7, 8, 9, 10, 11]. For $q\bar{q} \rightarrow t\bar{t}$ the soft anomalous dimension is a 2×2 matrix while that for $gg \rightarrow t\bar{t}$ is a 3×3 matrix.

The soft anomalous dimension matrix for $q\bar{q} \rightarrow t\bar{t}$ is

$$\Gamma_{Sq\bar{q}} = \begin{bmatrix} \Gamma_{q\bar{q}11} & \Gamma_{q\bar{q}12} \\ \Gamma_{q\bar{q}21} & \Gamma_{q\bar{q}22} \end{bmatrix}$$

At one loop, in a color basis of singlet and octet exchange in the s channel, we find [1, 11]

$$\begin{aligned}\Gamma_{q\bar{q}11}^{(1)} &= -C_F [L_\beta + 1], & \Gamma_{q\bar{q}12}^{(1)} &= \frac{C_F}{C_A} \ln \left(\frac{u_1}{t_1} \right), & \Gamma_{q\bar{q}21}^{(1)} &= 2 \ln \left(\frac{u_1}{t_1} \right), \\ \Gamma_{q\bar{q}22}^{(1)} &= C_F \left[4 \ln \left(\frac{u_1}{t_1} \right) - L_\beta - 1 \right] + \frac{C_A}{2} \left[-3 \ln \left(\frac{u_1}{t_1} \right) + \ln \left(\frac{t_1 u_1}{s m^2} \right) + L_\beta \right],\end{aligned}$$

where $L_\beta = [(1+\beta^2)/(2\beta)] \ln[(1-\beta)/(1+\beta)]$ with $\beta = \sqrt{1-4m^2/s}$.

In presenting our two-loop results, we use the short-hand notation M_β to denote the terms inside the curly brackets in our previous expression for the cusp anomalous dimension. Then at two loops we find [11]

$$\begin{aligned}\Gamma_{q\bar{q}11}^{(2)} &= \frac{K}{2} \Gamma_{q\bar{q}11}^{(1)} + C_F C_A M_\beta, & \Gamma_{q\bar{q}12}^{(2)} &= \frac{K}{2} \Gamma_{q\bar{q}12}^{(1)} - \frac{C_F}{2} N_\beta \ln \left(\frac{u_1}{t_1} \right), \\ \Gamma_{q\bar{q}21}^{(2)} &= \frac{K}{2} \Gamma_{q\bar{q}21}^{(1)} + C_A N_\beta \ln \left(\frac{u_1}{t_1} \right), & \Gamma_{q\bar{q}22}^{(2)} &= \frac{K}{2} \Gamma_{q\bar{q}22}^{(1)} + C_A \left(C_F - \frac{C_A}{2} \right) M_\beta,\end{aligned}$$

with N_β a subset of the terms of M_β .

For the $g(p_a) + g(p_b) \rightarrow t(p_1) + \bar{t}(p_2)$ channel we choose the color basis: $c_1 = \delta^{ab} \delta_{12}$, $c_2 = d^{abc} T_{12}^c$, $c_3 = i f^{abc} T_{12}^c$. The soft anomalous dimension matrix for $gg \rightarrow t\bar{t}$ in this basis is

$$\Gamma_{Sgg} = \begin{bmatrix} \Gamma_{gg11} & 0 & \Gamma_{gg13} \\ 0 & \Gamma_{gg22} & \Gamma_{gg23} \\ \Gamma_{gg31} & \Gamma_{gg32} & \Gamma_{gg33} \end{bmatrix}.$$

At one loop we have [1, 11]

$$\begin{aligned} \Gamma_{gg11}^{(1)} &= -C_F [L_\beta + 1], & \Gamma_{gg22}^{(1)} &= -C_F [L_\beta + 1] + \frac{C_A}{2} \left[\ln \left(\frac{t_1 u_1}{m_s^2} \right) + L_\beta \right], & \Gamma_{gg23}^{(1)} &= \frac{C_A}{2} \ln \left(\frac{u_1}{t_1} \right), \\ \Gamma_{gg13}^{(1)} &= \ln \left(\frac{u_1}{t_1} \right), & \Gamma_{gg31}^{(1)} &= 2 \ln \left(\frac{u_1}{t_1} \right), & \Gamma_{gg32}^{(1)} &= \frac{N_c^2 - 4}{2N_c} \ln \left(\frac{u_1}{t_1} \right). \end{aligned}$$

At two loops we find [11]

$$\begin{aligned} \Gamma_{gg11}^{(2)} &= \frac{K}{2} \Gamma_{gg11}^{(1)} + C_F C_A M_\beta, & \Gamma_{gg22}^{(2)} &= \frac{K}{2} \Gamma_{gg22}^{(1)} + C_A \left(C_F - \frac{C_A}{2} \right) M_\beta, & \Gamma_{gg23}^{(2)} &= \frac{K}{2} \Gamma_{gg23}^{(1)}, \\ \Gamma_{gg13}^{(2)} &= \frac{K}{2} \Gamma_{gg13}^{(1)} - \frac{C_A}{2} N_\beta \ln \left(\frac{u_1}{t_1} \right), & \Gamma_{gg31}^{(2)} &= \frac{K}{2} \Gamma_{gg31}^{(1)} + C_A N_\beta \ln \left(\frac{u_1}{t_1} \right), & \Gamma_{gg32}^{(2)} &= \frac{K}{2} \Gamma_{gg32}^{(1)}. \end{aligned}$$

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